

University of Illinois at Chicago
Department of Physics

Quantum Mechanics
Qualifying Examination

January 4, 2010.
9:00 am – 12:00 pm

Full credit can be achieved from completely correct answers to 4 questions. If the student attempts all 5 questions, all of the answers will be graded, and the top 4 scores will be counted toward the exam's total score.

Various equations, standard integrals, etc. are provided on the last page of the exam.

Question 1

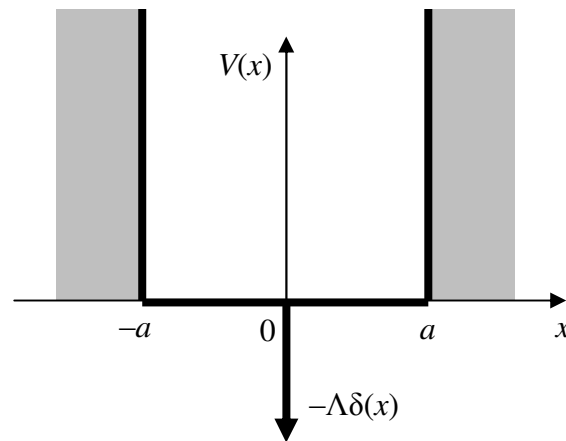
In heavy (large Z) hydrogen-like atoms, where to a very good approximation the reduced mass μ is equal to the electron mass m_e , relativistic corrections to the electron's kinetic energy need to be taken into account due to its large orbital velocity.

- a) Show that the first-order correction to the Hamiltonian due to the electron's relativistic kinetic energy is given by $\hat{H}_1 = -\frac{1}{2m_e c^2} \left(\frac{\mathbf{p}^2}{2m_e} \right)^2$.
- b) Verify that $\left\langle \frac{1}{r} \right\rangle = \frac{Z}{a_0}$ for the $1s$ state of the hydrogen-like atom.
- c) Similarly, evaluate $\left\langle \frac{1}{r^2} \right\rangle$ for the $1s$ state of the hydrogen-like atom.
- d) Use the fact that the unperturbed Hamiltonian is $\hat{H}_0 = \frac{\mathbf{p}^2}{2m_e} - \frac{Ze^2}{4\pi\epsilon_0 r}$, and the results from parts (a), (b) and (c), to evaluate the first-order perturbation to the electron's kinetic energy due to the relativistic correction.

Question 2

Consider an *attractive* delta-function potential $-\Lambda\delta(x)$, where the parameter Λ characterizes the strength of the delta-function, positioned at the center of an infinite square well of width $2a$; that is, a potential given by

$$\begin{aligned} V(x \leq -a) &= \infty \\ V(-a < x < 0) &= 0 \\ V(x = 0) &= -\Lambda\delta(x) \\ V(0 < x < a) &= 0 \\ V(x \geq a) &= \infty \end{aligned}$$



- a) Show that the equation determining the energy E of the eigenstate bound to the delta-function is given by

$$\tanh(\kappa a) = \frac{\hbar^2 \kappa}{m\Lambda},$$

where $\hbar\kappa = \sqrt{2m|E|}$ and m is the mass of the particle.

- b) What is the minimum strength of the delta-function for which a state with $E < 0$ exists?

Question 3

Quantum dots are important material systems in nanoscience.

Consider an electron of charge e and mass m_e confined in an idealized quantum dot with $V(r) = 0$ for $r < a$ and $V(r) = \infty$ for $r > a$; that is, an infinite spherical well of radius a . The eigenstates of this potential are given by

$$\langle \mathbf{r} | \Psi_{nlm} \rangle = C_{nl} j_l \left(\frac{s_{nl} r}{a} \right) Y_{lm}(\theta, \phi),$$

where $Y_{lm}(\theta, \phi)$ is the spherical harmonic associated with the usual angular momentum quantum numbers l and m , and C_{nl} is the normalization constant for the l^{th} spherical Bessel function j_l with roots s_{nl} . The first three spherical Bessel functions with $\rho = s_{nl} r/a$ are

$$j_0(\rho) = \frac{\sin \rho}{\rho} \quad j_1(\rho) = \frac{\sin \rho}{\rho^2} - \frac{\cos \rho}{\rho} \quad j_2(\rho) = \left(\frac{3}{\rho^3} - \frac{1}{\rho} \right) \sin \rho - \frac{3}{\rho^2} \cos \rho$$

and the first few roots are given below:

s_{nl}	$l = 0$	$l = 1$	$l = 2$	$l = 3$	$l = 4$
$n = 1$	π	4.49	5.76	6.99	9.36
$n = 2$	2π	7.73	9.10	10.42	
$n = 3$	3π				

The spherical Bessel functions also satisfy the relation

$$\int_0^a r^2 dr j_l \left(\frac{s_{nl} r}{a} \right) j_l \left(\frac{s_{n'l'} r}{a} \right) = \frac{a^3}{2} [j_{l-1}(s_{nl})]^2 \delta_{nn'}.$$

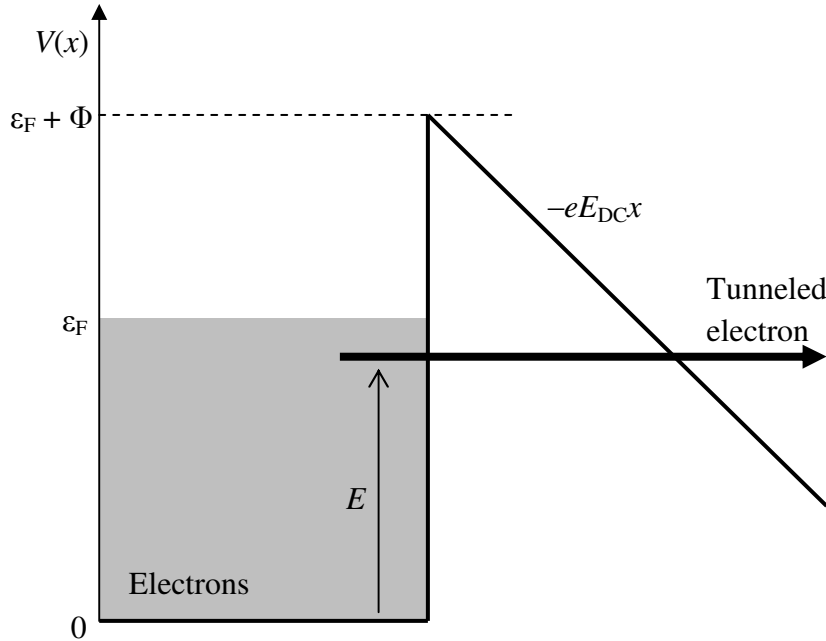
- What are the energies of the lowest four eigenstates?
- Evaluate the normalization constants C_{10} and C_{11} .
- Consider the state $\langle \mathbf{r} | \psi \rangle = A \left[\left(\frac{\sin \rho}{\rho^2} - \frac{\cos \rho}{\rho} \right) \sin \theta \cos \phi + \frac{\sin \rho}{\rho} \right]$, where A is a constant.

With what probability is the ground state energy measured? And what is the probability of measuring a z -component of the angular momentum with a value $+\hbar$?

- In the presence of a strong homogeneous magnetic field \mathbf{B} , the energy states split – the Zeeman effect. Neglecting spin (the normal Zeeman effect), the additional contribution to the Hamiltonian can be written as $H_1 = \frac{e}{2m_e} \mathbf{L} \cdot \mathbf{B}$. Amongst the four lowest states identified in part (a), at what magnetic field strength will states originating from one energy level overlap energetically with those from another?

Question 4

Field emission, a quantum tunneling phenomenon, occurs when a DC electric field E_{DC} applied perpendicular to a material surface becomes sufficiently strong to allow electrons in the material to tunnel out into the vacuum.



In materials like metals, the zero-point energy E of the electrons in the metal is usually defined in terms of the Fermi energy ε_F (the energy of the last filled electronic state at zero Kelvin), which is located at an energy Φ (the material work function) below the vacuum energy.

- a) Using the approximation for the one-dimensional tunneling probability through a ‘thick’ barrier,

$$|T|^2 \approx \exp \left[-2 \int_a^b dx \sqrt{\frac{2m}{\hbar^2} (V(x) - E)} \right],$$

of thickness $b - a$, determine the transmission probability for an electron at an arbitrary energy E ; that is, at an energy $E' = \varepsilon_F + \Phi - E$ below the top of the barrier.

- b) In a real material at non-zero temperature T , the energy distribution of electrons is given by the Fermi function

$$f(E) = \frac{1}{1 + \exp \left(-\frac{\varepsilon_F - E}{k_B T} \right)},$$

where k_B is Boltzmann’s constant. For $\exp[(\Phi - E')/k_B T] \gg 1$, determine the energy below the top of the barrier for which electron field emission is most probable.

Question 5

Consider the following *spatial-coordinate-independent* two-particle Hamiltonian for a spin system consisting of two *identical* spin $\frac{1}{2}$ particles:

$$\hat{H} = A + B\mathbf{S}_1 \cdot \mathbf{S}_2$$

where A and B are constants.

- a) Find the 4×4 matrix representation of \hat{H} in the uncoupled basis set

$$|\uparrow\uparrow\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}, |\uparrow\downarrow\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}, |\downarrow\uparrow\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}, \text{ and } |\downarrow\downarrow\rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix},$$

where $|\uparrow\rangle$ represents a spin-up particle with $S_z = \frac{1}{2}\hbar$ and a spin-down particle with $S_z = -\frac{1}{2}\hbar$ is $|\downarrow\rangle$.

- b) Calculate the eigenvalues of \hat{H} .
c) What are the eigenstates of the diagonalized Hamiltonian?
d) Verify your results by finding the eigenvalues of the four states in the coupled basis $|S, M\rangle$ representing the total spin S and its z -component M .

Equation Sheet

$$Y_{00}(\theta, \phi) = \frac{1}{\sqrt{4\pi}} \quad Y_{10}(\theta, \phi) = \sqrt{\frac{3}{4\pi}} \cos \theta \quad Y_{1\pm 1}(\theta, \phi) = \mp \sqrt{\frac{3}{8\pi}} e^{\pm i\phi} \sin \theta$$

$$Y_{20}(\theta, \phi) = \sqrt{\frac{5}{16\pi}} (3 \cos^2 \theta - 1) \quad Y_{2\pm 1}(\theta, \phi) = \mp \sqrt{\frac{15}{8\pi}} e^{\pm i\phi} \sin \theta \cos \theta$$

$$Y_{2\pm 2}(\theta, \phi) = \sqrt{\frac{15}{32\pi}} e^{\pm 2i\phi} \sin^2 \theta$$

$$E_n = - \left[\frac{\mu}{2\hbar^2} \left(\frac{Ze^2}{4\pi\epsilon_0} \right)^2 \right] \frac{1}{n^2} = -13.6 \frac{Z}{n^2} \text{ eV} \quad a_0 = \frac{4\pi\epsilon_0 \hbar^2}{\mu e^2} \quad \frac{1}{\mu} = \frac{1}{m_e} + \frac{1}{M_{\text{nucleus}}}$$

$$R_{10}(r) = 2 \left(\frac{Z}{a_0} \right)^2 \exp \left[-\frac{Zr}{a_0} \right] \quad R_{20}(r) = 2 \left(\frac{Z}{2a_0} \right)^2 \left(1 - \frac{Zr}{2a_0} \right) \exp \left[-\frac{Zr}{2a_0} \right]$$

$$R_{21}(r) = \frac{1}{\sqrt{3}} \left(\frac{Z}{2a_0} \right)^2 \frac{Zr}{a_0} \exp \left[-\frac{Zr}{2a_0} \right]$$

$$\int_0^{\infty} dx x^m \exp(-ax^2) = \frac{\Gamma\left(\frac{m+1}{2}\right)}{2a^{(m+1)/2}} \quad ; \quad \Gamma(n+1) = n\Gamma(n), \Gamma(n+1) = n!, \Gamma\left(\frac{1}{2}\right) = \sqrt{\pi}$$

$$\text{so that } \int_0^{\infty} dx x^{2n} \exp(-\lambda^2 x^2) = \frac{1.3.5 \dots (2n+1)\sqrt{\pi}}{2^n \lambda^{2n+1}}$$

$$\int_0^{\infty} dx x^n e^{-\lambda x} = \frac{n!}{\lambda^{n+1}}$$

$$\int dx \sqrt{A+Bx} = \frac{2}{3B} (A+Bx)^{3/2} \quad \int dx x \sqrt{A+Bx} = -\frac{2(2A-3Bx)(A+Bx)^{3/2}}{15B^2}$$